

## Isomerism of Benzoquinone Monoximes (Nitrosophenols). XII. The Crystal Structure of $\beta$ -5-n-Propoxy-*o*-quinone-2-oxime\*

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The crystal structure of  $\beta$ -5-n-propoxy-*o*-quinone-2-oxime has been determined by comparison with those of the isomorphous  $\beta$ -5-(2'-bromoethoxy)- and  $\beta$ -5-(2'-chloroethoxy)-*o*-quinone-2-oxime. The unit-cell dimensions of the propoxy compound at  $-120^\circ\text{C}$  are  $a=10.64$ ,  $b=12.57$ ,  $c=6.63$  Å,  $\beta=95^\circ 31'$ ,  $Z=4$ . The space group is  $P2_1/c$ . The structure was refined by the method of least-squares with data collected at  $-120^\circ\text{C}$ . The mean value of two short distances (1.35 Å) and four long distances (1.45 Å) in the hexagonal ring and the length of the carboxyl bond (1.27 Å) indicate a predominantly oxime-like structure. The intramolecular hydrogen-bonded NOH and CO groups are in the *syn* position. The  $\text{CH}_3\text{-C}_2\text{H}_4\text{-O}$ -group has the *gauche* conformation. The molecule is planar with the exception of the methyl group which is 1.22 Å outside the plane of the molecule. Exceptionally short intermolecular distances (3.12, 3.25 Å) are found between molecules in adjacent layers in the direction of the *c* axis.

### Introduction

As a part of the investigations on the tautomerism of quinone-oximes (nitrosophenols) in this laboratory Umans (1959) studied the properties of 5-alkoxy-*o*-quinone-2-oximes which occur in two forms. It was concluded by Romers & Umans (1960), in a preliminary account of the crystal structures of some of these oximes, that the stable orange  $\alpha$  forms possess the oxime structure. Bartindale, Crowder & Morley (1959) arrived at the same conclusion in their investigation of the crystal structure of  $\alpha$ -5-methoxy-*o*-quinone-2-oxime. The first mentioned authors suspected that the green  $\beta$  forms might have a nitrosophenolic character. This hypothesis was mainly supported by the optical properties of the crystals, because the two-dimensional X-ray analysis was not sufficiently accurate to distinguish between the two forms.

Not all  $\beta$  forms are unstable. Umans (1959) proved that the green  $\beta$  form of 5-n-propoxy-*o*-quinone-2-oxime ( $\beta$ -PQO) is stable under normal conditions, while it is difficult to obtain its corresponding orange  $\alpha$  form. Since, moreover,  $\beta$ -PQO is isomorphous with  $\beta$ -5-(2'-bromoethoxy) ( $\beta$ -Br EQO) and  $\beta$ -5-(2'-chloroethoxy)-*o*-quinone-2-oxime ( $\beta$ -Cl EQO), the first mentioned compound is ideally suited for a three-dimensional X-ray diffraction analysis.

### Morphology and optics

Crystals of  $\beta$ -PQO, melting point  $94\text{--}95^\circ\text{C}$ , are lath-shaped prisms with well developed faces (100) and elongated in the direction [001]. They can be cleaved parallel to [010] and [001]. However, persistent

twinning with (100) as common twinning plane was found after cleavage parallel to [010].

The crystals have a dark metallic lustre and exhibit a strongly marked pleochroism on inspection in polarized light. When observed through the face (100) with [001] parallel to the direction of the electric vector of polarized light their colour is yellow-green. The colour changes into orange on turning the crystals  $90^\circ$  with respect to the direction of the electric vector. This pleochroism is even more pronounced in the corresponding chloro compound  $\beta$ -Cl EQO which shows a variation between grass-green and garnet-orange. The absorption spectrum of a suspension of  $\beta$ -PQO in Nujol shows maxima at  $653\text{ m}\mu$  and  $323\text{ m}\mu$ ; its solution in benzene has corresponding maxima at  $645\text{ m}\mu$  (molecular extinction coefficient  $\epsilon=35$ ) and at  $325\text{ m}\mu$  ( $\epsilon=11000$ ).

The crystals are biaxially negatively birefringent ( $n_\alpha \ll n_\beta < n_\gamma$ ). The obtuse bisectrix interference figure is observable through (100) in convergent polarized light. The angle  $2V$  is close to  $90^\circ$  and the difference  $(n_\gamma - n_\beta) - (n_\beta - n_\alpha)$  is small.  $n_\beta$  is along [010] and  $n_\alpha$  and  $n_\gamma$  are approximately in the respective directions [001] and [100].

### Experimental

Crystals of PQO with cross sections of approximately  $0.3 \times 0.3$  mm were cut in the directions [001] and [010]. The cell dimensions at  $20^\circ\text{C}$  and at  $-120^\circ\text{C}$  were obtained from zero-layer Weissenberg photographs about [010] and [001] (copper radiation,  $\lambda=1.5418$  Å) superposed with aluminum powder lines for calibration purposes. The cell dimensions of  $\beta$ -Br EQO,  $\beta$ -Cl EQO and  $\beta$ -PQO are summarized in Table 1.

Their space groups are  $P2_1/c$  and there are four molecules per unit cell. The visually estimated re-

\* See part VIII (Romers & Fischman, 1960) for earlier papers of this series.

Table 1. *Dimensions* (Å) *and densities* (g.cm<sup>-3</sup>) *of some*  $\beta$ -5-alkoxy-o-quinone-2-oximes

Compound	<i>a</i>	<i>b</i>	<i>c</i>	$\beta$	<i>D<sub>m</sub></i>	<i>D<sub>x</sub></i>	Temperature
$\beta$ -Br EQO	10.86	12.73	6.76	90	1.77	1.75	20 °C
$\beta$ -Cl EQO	10.70	12.28	6.59	90	1.48	1.54	-120 °C
$\beta$ -PQO	10.66 ± 0.01	12.67 ± 0.01	6.79 ± 0.02	95° 36' ± 20'	1.29	1.31	20 °C
$\beta$ -PQO	10.64 ± 0.01	12.57 ± 0.02	6.63 ± 0.01	96° 31' ± 30'			-120 °C

flexion intensities were recorded at -120 °C on multiple-film equi-inclination Weissenberg photographs with Cu *K* $\alpha$  radiation. The crystals were dipped into liquid air in order to minimize extinction effects. The first six layers about [001] and the first nine layers about [010] were prepared according to a low temperature technique developed in this laboratory by Altona & Romers (1963). The maximum number of independent reflexions within the copper radiation sphere is 2123. Of these, 1613 reflexions could have been recorded, but only 1332 were observed. The usual Lorentz and polarization factors and film spot corrections were applied. The intensities were not corrected for absorption, because the linear absorption coefficient for copper radiation is only 9.8 cm<sup>-1</sup>.

#### Determination of the structure

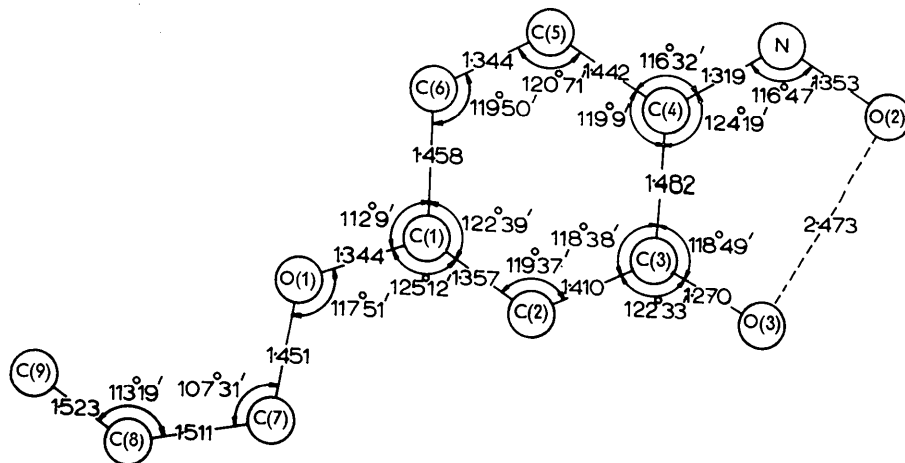
Crystals of  $\beta$ -Br EQO were first obtained. Inspection of their optical properties and unit-cell dimensions indicated that the [001] projection was the most favourable for investigation. However, the *x* and *y* parameters of the bromine atom were the sole information which could be extracted from the Patterson function *P*(*uv*), because the application of minimum functions (Buerger, 1959) failed to reveal the structure. The electron density projection  $\rho(xy)$  based upon the phases of the bromine atom contributions to the structure factors could not be interpreted in terms of a reasonable molecular model. Afterwards we realized that the failure of the vector shift method as well as of the heavy atom technique was due to the poor

quality of the diffraction data (absorption errors and twinning) and to the ambiguity of the projection, which shows several possible hexagonal arrays of peaks (Romers & Umans, 1960, Figs. 2 and 3; see also Fig. 4 of this paper).

Meanwhile crystals of  $\beta$ -Cl EQO became available. This compound changes spontaneously at room temperature into the stable orange  $\alpha$  form. Moreover, the change is accelerated by X-rays. We only succeeded in collecting *hk0* and *hk1* reflexion data at -120 °C. A comparison of the corresponding reflexions *hk0* of  $\beta$ -Br EQO and  $\beta$ -Cl EQO indicated that their structures are isomorphous. The signs of most structure factors *F*(*hk0*) of the chloro compound were determined and its structure was found by calculation of generalized electron density functions  $\rho_0(xy)$ , *C*<sub>1</sub>(*xy*) and *S*<sub>1</sub>(*xy*) (Romers & Umans, 1960, Figs. 2 and 3). The function  $\rho_0(xy)$  is the ordinary projection of the electron density; the latter two functions are the cosine function *C*<sub>1</sub> and the sine function *S*<sub>1</sub> based upon two-dimensional Fourier summations with Fourier coefficients *F*(*hk1*) (Cochran & Dyer, 1952; Romers & Fischmann, 1960; Fischmann, MacGillavry & Romers, 1961).

#### The refinement

The remarkable stability of  $\beta$ -PQO led to the measurement of its unit cell and to the fortuitous discovery that its crystals are isomorphous with the above-mentioned corresponding halogeno compounds. The three-dimensional refinement of the structure of  $\beta$ -PQO by the method of least-squares was based upon

Fig. 1. Intramolecular bond distances (Å) and bond angles in  $\beta$ -PQO.

the positional parameters obtained from the diffraction analysis of the chloro compound. The numbering of the atoms is indicated in Fig. 1. The atomic scattering factors of carbon, nitrogen and oxygen used in the structure factor calculations were those of Berghuis, Haanappel, Potters, Loopstra, MacGillavry & Veenendaal (1955). McWeeny's (1951) form factors for hydrogen were used in the final stages of refinement.

After the first three cycles, in which one overall isotropic temperature factor was applied, the reliability index  $R = \Sigma ||F_o| - |F_c|| / \Sigma |F_o|$  dropped from 28% to 22%. These calculations were performed on ZEBRA computers at Utrecht and Groningen and a program with diagonal approximation, written by Dr J. C. Schoone, was employed.

The calculations were continued on the Mercury computer at Manchester with least-squares programs written by Mr O. S. Mills (Mills & Rollett, 1961). The following weighting scheme was used:  $w = 1/(1 + F_o^2/8F_{\min}^2)$ , where  $|F_{\min}|$  is the observable minimum of  $|F_o|$ . The next three cycles were carried out with isotropic individual temperature factors. At this stage the reliability index  $R$  was 16.7%. A difference map  $\rho_o - \rho_c$  was calculated. Important positional shifts of the heavy atoms were not found, but the map gave indications of the vibrational movements of the heavy atoms and showed distinctly the positions of the hydrogen atoms.

The following two cycles (Nos. 7 and 8) were calculated with anisotropic temperature factors. In view of the small positional shifts after the eighth cycle ( $< 0.0013 \text{ \AA}$ ) the refinement could be considered completed. However, the reliability index was still rather high (14.8%) and a careful survey of all structure factors revealed the poor quality of some non-zero Weissenberg photographs. Moreover, it had been discovered that the formerly measured length of the  $c$  axis (Romers & Umans, 1961) was  $0.18 \text{ \AA}$  too long.

New diffraction data were collected at  $-120 \text{ }^\circ\text{C}$  for the faulty levels, resulting in a reliability index of 11.84%. Hydrogen atoms were placed at the positions indicated by the difference Fourier map, but their contributions to the structure factors reduced  $R$  only to 11.65%. The hydrogen atoms contribute

mainly to the very strong reflexions 002 and 102 which define the plane of the molecule and the reflexion 100. But these three reflexions suffer seriously from extinction, which explains the negligible influence of the hydrogen atoms on the reliability index.

Two more cycles of least-squares (Nos. 9 and 10) were calculated with the hydrogen contributions added to the structure factors. They were kept in fixed positions and their scattering factors (McWeeny, 1951) were multiplied with a constant overall isotropic temperature factor  $B = 2.22 \text{ \AA}^{-2}$ . The last cycle but one gave three shifts:  $\Delta x(\text{C}(1)) = -0.016$ ,  $\Delta y(\text{C}(2)) = 0.011$  and  $\Delta z(\text{C}(9)) = -0.015 \text{ \AA}$ ; all other positional shifts were much smaller. The last cycle gave only minor shifts (average value  $0.0016 \text{ \AA}$ ) with the exception of  $\Delta x(\text{C}(1)) = +0.0057$  and  $\Delta z(\text{C}(9)) = +0.0092 \text{ \AA}$ . In view of these oscillations it might have been better to continue the refining process with another cycle with reduced shifts for the atoms C(1) and C(9), but this intention was prevented by lack of computing facilities. However, the reliability

Table 2. *Final positional parameters in fractions of cell edges*

	$x$	$y$	$z$
C(1)	0.2845	0.0241	0.3010
C(2)	0.3864	0.9640	0.2672
C(3)	0.5026	0.0140	0.2431
C(4)	0.5101	0.1315	0.2575
C(5)	0.3988	0.1909	0.2952
C(6)	0.2900	0.1398	0.3170
C(7)	0.1486	0.8723	0.3023
C(8)	0.0139	0.8497	0.3431
C(9)	0.9167	0.9034	0.1928
N	0.6133	0.1873	0.2383
O(1)	0.1686	0.9861	0.3233
O(2)	0.7185	0.1318	0.2061
O(3)	0.6012	0.9622	0.2109
H(1)	0.929	0.872	0.062
H(2)	0.815	0.887	0.200
H(3)	0.928	0.976	0.153
H(4)	0.997	0.890	0.456
H(5)	0.005	0.230	0.357
H(6)	0.200	0.842	0.423
H(7)	0.155	0.848	0.154
H(8)	0.392	0.888	0.267
H(9)	0.405	0.270	0.316
H(10)	0.198	0.163	0.333
H(11)	0.694	0.065	0.267

Table 3. *Vibrational parameters*

	$b_{11}$	$b_{12}$	$b_{13}$	$b_{22}$	$b_{23}$	$b_{33}$
C(1)	0.00446	0.00054	0.00239	0.00276	0.00026	0.01228
C(2)	0.00440	0.00085	0.00098	0.00268	0.00173	0.01230
C(3)	0.00465	0.00060	-0.00006	0.00307	-0.00067	0.01086
C(4)	0.00448	-0.00084	-0.00108	0.00285	0.00109	0.01265
C(5)	0.00610	0.00059	0.00286	0.00231	-0.00011	0.01586
C(6)	0.00536	0.00086	0.00438	0.00262	-0.00068	0.01832
C(7)	0.00461	-0.00089	0.00259	0.00256	-0.00076	0.01926
C(8)	0.00467	-0.00156	0.00754	0.00459	0.00011	0.02303
C(9)	0.00413	-0.00014	0.00021	0.00542	-0.00253	0.03209
N	0.00487	-0.00042	0.00138	0.00322	-0.00036	0.01722
O(1)	0.00369	0.00057	0.00424	0.00271	-0.00053	0.02019
O(2)	0.00439	-0.00145	0.00286	0.00349	-0.00023	0.02501
O(3)	0.00397	0.00014	0.00288	0.00308	-0.00091	0.01649



Table 4 (cont.)

FC	62	63	64	65	66	67	68	69	70	71	72	73	74	75	76	77	78	79	80	81	82	83	84	85	86	87	88	89	90	91	92	93	94	95	96	97	98	99	100		
FO	110	111	112	113	114	115	116	117	118	119	120	121	122	123	124	125	126	127	128	129	130	131	132	133	134	135	136	137	138	139	140	141	142	143	144	145	146	147	148	149	150
FKI	150	151	152	153	154	155	156	157	158	159	160	161	162	163	164	165	166	167	168	169	170	171	172	173	174	175	176	177	178	179	180	181	182	183	184	185	186	187	188	189	190
FC	190	191	192	193	194	195	196	197	198	199	200	201	202	203	204	205	206	207	208	209	210	211	212	213	214	215	216	217	218	219	220	221	222	223	224	225	226	227	228	229	230
FO	230	231	232	233	234	235	236	237	238	239	240	241	242	243	244	245	246	247	248	249	250	251	252	253	254	255	256	257	258	259	260	261	262	263	264	265	266	267	268	269	270
FKI	270	271	272	273	274	275	276	277	278	279	280	281	282	283	284	285	286	287	288	289	290	291	292	293	294	295	296	297	298	299	300	301	302	303	304	305	306	307	308	309	310
FC	310	311	312	313	314	315	316	317	318	319	320	321	322	323	324	325	326	327	328	329	330	331	332	333	334	335	336	337	338	339	340	341	342	343	344	345	346	347	348	349	350
FO	350	351	352	353	354	355	356	357	358	359	360	361	362	363	364	365	366	367	368	369	370	371	372	373	374	375	376	377	378	379	380	381	382	383	384	385	386	387	388	389	390
FKI	390	391	392	393	394	395	396	397	398	399	400	401	402	403	404	405	406	407	408	409	410	411	412	413	414	415	416	417	418	419	420	421	422	423	424	425	426	427	428	429	430
FC	430	431	432	433	434	435	436	437	438	439	440	441	442	443	444	445	446	447	448	449	450	451	452	453	454	455	456	457	458	459	460	461	462	463	464	465	466	467	468	469	470
FO	470	471	472	473	474	475	476	477	478	479	480	481	482	483	484	485	486	487	488	489	490	491	492	493	494	495	496	497	498	499	500	501	502	503	504	505	506	507	508	509	510
FKI	510	511	512	513	514	515	516	517	518	519	520	521	522	523	524	525	526	527	528	529	530	531	532	533	534	535	536	537	538	539	540	541	542	543	544	545	546	547	548	549	550
FC	550	551	552	553	554	555	556	557	558	559	560	561	562	563	564	565	566	567	568	569	570	571	572	573	574	575	576	577	578	579	580	581	582	583	584	585	586	587	588	589	590
FO	590	591	592	593	594	595	596	597	598	599	600	601	602	603	604	605	606	607	608	609	610	611	612	613	614	615	616	617	618	619	620	621	622	623	624	625	626	627	628	629	630
FKI	630	631	632	633	634	635	636	637	638	639	640	641	642	643	644	645	646	647	648	649	650	651	652	653	654	655	656	657	658	659	660	661	662	663	664	665	666	667	668	669	670
FC	670	671	672	673	674	675	676	677	678	679	680	681	682	683	684	685	686	687	688	689	690	691	692	693	694	695	696	697	698	699	700	701	702	703	704	705	706	707	708	709	710
FO	710	711	712	713	714	715	716	717	718	719	720	721	722	723	724	725	726	727	728	729	730	731	732	733	734	735	736	737	738	739	740	741	742	743	744	745	746	747	748	749	750
FKI	750	751	752	753	754	755	756	757	758	759	760	761	762	763	764	765	766	767	768	769	770	771	772	773	774	775	776	777	778	779	780	781	782	783	784	785	786	787	788	789	790

index was reduced only from 11.65% to 11.60%, indicating that convergence had been attained. The final calculated structure factors from cycle 10 are listed in Table 4.

#### The parameters and the accuracy

The final positional parameters in fractions of the cell edges are listed in Table 2. The mean positional standard deviation (0.0035 Å) obtained from inverted matrix diagonal elements in cycle 10 is surprisingly low. Its value is an underestimate since the absolute weights were unknown during the refinement and the adopted weight scheme may have obscured the influence of very weak and very strong reflexions. Therefore, the mean positional standard deviation was estimated by the method of Cruickshank (1960). The standard deviation found in this way (0.0055 Å) is, according to Cruickshank, perhaps 20% too high, but it seems safe to accept this value. The standard deviations in the cell edges are approximately 0.002 Å per measured length of 1.5 Å. It follows that the mean standard deviations in the interatomic bond lengths are  $(2 \times 0.0055^2 + 0.002^2)^{1/2} = 0.008$  Å, giving a 1% significance level of 0.024 Å.

The final vibrational parameters, defining a temperature factor

$$\exp(-b_{11}h^2 - b_{12}hk - b_{13}hl - b_{22}k^2 - b_{23}kl - b_{33}l^2),$$

are listed in Table 3. It has been shown (Altona, Knobler & Romers, 1963) that the absolute values of the  $b$  parameters are doubtful in the case of visual estimation of photographically recorded intensities. The question arises whether the data are sufficiently accurate to allow for anisotropic refinement. Since no program was immediately available to determine the principal axes of the ellipsoid of vibration of all atoms, we confined ourselves to calculate the *mean* principal vibration directions of the eight atoms C(1) ··· C(6), N and O(3), belonging to the planar part of the molecule (see next section). A similar procedure has been applied by Stam & MacGillavry (1963) according to a method proposed by Waser (1955) and improved by Busing & Levy (1958).

The largest vibration with mean square displacement  $\langle u_1^2 \rangle = 0.0344$  Å<sup>2</sup>, or  $B_1 = 2.72$  Å<sup>2</sup>, makes an angle of 23° with the normal to the above-mentioned plane of the molecule. The second largest vibration with  $\langle u_2^2 \rangle = 0.0253$  Å<sup>2</sup>, or  $B_2 = 1.99$  Å<sup>2</sup>, makes an angle of 44° with the normal and of 49° with the positive  $a$  axis. The smallest vibration  $\langle u_3^2 \rangle = 0.0214$  Å<sup>2</sup>, or  $B_3 = 1.69$  Å<sup>2</sup> makes an angle of 81° with the normal and of -57° with the positive  $a$  axis. These numbers roughly agree with the indications of the difference Fourier synthesis made during an earlier stage of the refinement. Judging from the rather small  $B$  values found in this determination we estimate that a

correction of at most 0.005 Å owing to the rotational oscillations of the molecules should be added to the distances involving the tail atoms C(7), C(8) and C(9). These corrections are at most 0.003 Å for all distances involving the remaining atoms of the molecule. In view of the uncertainties of the absolute values of the *b* parameters and of the lack of computational facilities these corrections have been omitted.

### Discussion of the structure

The distances and bond angles between the atoms in the structure were calculated on the X1 computer of the University of Leiden with the aid of a program written by Geise, Jonkers & Romers (1963). The intramolecular bond distances and bond angles are indicated in Fig. 1. The ring atoms C(1)···C(6) and the atoms N and O(3) define a plane

$$0.1598X - 0.0726Y + Z - 2.4213 = 0$$

with respect to an orthogonal coordinate system *XYZ*, where *X* and *Y* have the directions [100] and [010] respectively, and *Z* is perpendicular to (001). The deviations of atoms C(1)···C(6) and atoms N and O(3) from this plane (Table 5) are within the standard deviations of the coordinates. The atoms O(1) and O(2) are possibly and C(7) and C(8) are significantly outside of this plane, C(9) being at a large distance (1.20 Å).

Table 5. Distances from the molecular plane

C(1)	-0.0058 Å	C(8)	-0.0337 Å
C(2)	+0.0012	C(9)	-1.2048
C(3)	-0.0030	N	-0.0025
C(4)	-0.0023	O(1)	-0.0227
C(5)	-0.0004	O(2)	+0.0168
C(6)	+0.0036	O(3)	+0.0019
C(7)	-0.0894		

The NO(H) group is *syn* with respect to the CO(H) group and the intramolecular distance O(2)–O(3) (2.473 Å) is comparable to the intramolecular hydrogen bonds found in salicylic acid (2.590 Å, Cochran, 1953), maleic acid (2.46 Å, Shahat, 1952), and cytosine-5-acetic acid (2.506 Å, Marsh, Bierstedt & Eichhorn,

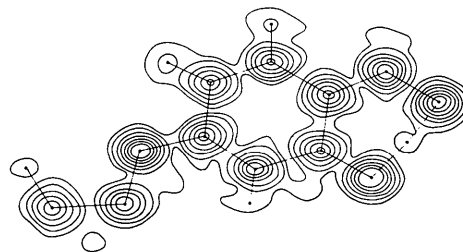


Fig. 2. Electron density map of  $\beta$ -PQO calculated in the plane of the molecule. The contour lines are drawn at 1, 2, 4, 6 ... e.Å<sup>-3</sup>.

1962). A Fourier synthesis was calculated (Fig. 2) in the plane of the molecule. This calculation was performed on the computer Pascal, Philips Laboratories, Eindhoven with a program devised by Dr P. Braun. The Fourier map clearly shows the hydrogen atoms bound to the carbon atoms C(2), C(5) and C(6). A fourth hydrogen atom is found half way between the oxygen atoms O(2) and O(3). Its peak value (1.5 e.Å<sup>-3</sup>) is, however, lower than those of the other hydrogen atoms.

The bond distances C(1)–C(2) (1.357 Å) and C(5)–C(6) (1.344 Å) in *ortho* positions are significantly shorter than the four other bond distances in the ring. The C(3)–O(3) bond (1.270 Å) and the C(4)–N bond (1.319 Å) have a pronounced double-bond character. The distribution of longer and shorter bond distances in the ring as well as the short C(3)–O(3) and C(4)–N bonds clearly indicate a predominantly oxime-like structure. The molecular structure might be described in terms of resonance between the structures (a), (b), (c) ... (h) of Fig. 3.

When these structures are given weights as indicated in Fig. 3 and when bond orders *p* and distances *R*<sub>1</sub> and *R*<sub>2</sub> are used as stated in Table 6 the bond lengths calculated from

$$R = R_1 - (R_1 - R_2) \times 1.333p / (0.333p + 1)$$

(Cruickshank & Sparks, 1960; Cruickshank, 1962) are as shown in Table 6. The standard deviation between the calculated and observed values is 0.015 Å.

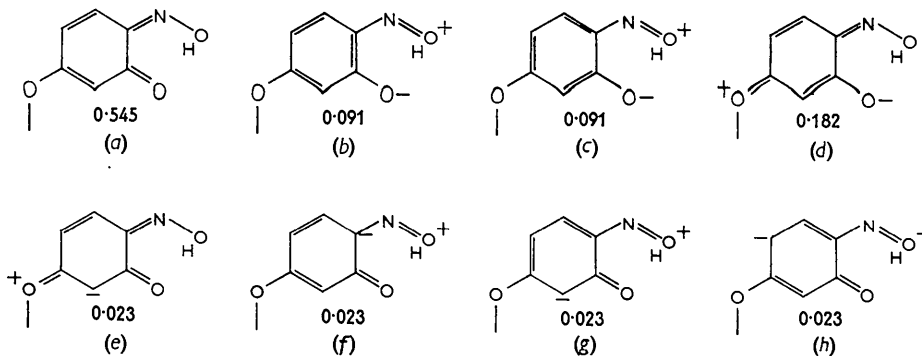


Fig. 3. Canonical 'structures' of  $\beta$ -PQO.



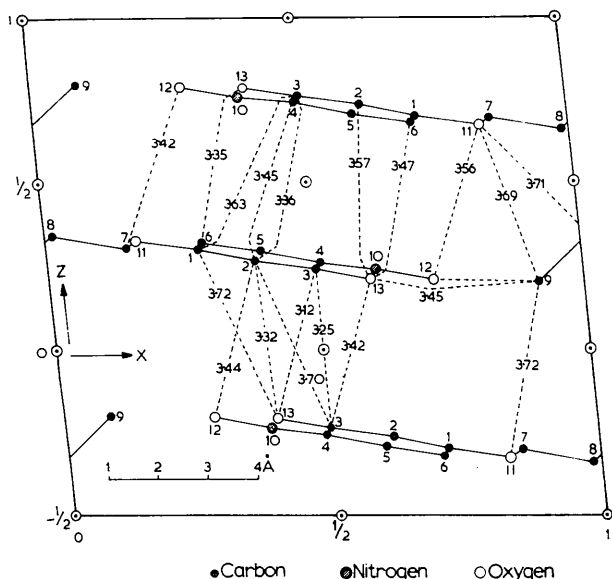


Fig. 5. Part of the crystal structure of  $\beta$ -PQO viewed along the  $b$  axis. The screw axis operation has been omitted.

even shorter distances between oxygen and carbon atoms have been found in the crystal structures of chloranil (Chu, Jeffrey & Sakurai, 1962) and of parabanic acid (Davies & Blum, 1955). The packing in these last-mentioned structures is of a herring-bone type;  $\beta$ -PQO packs more or less in the form of a layer structure. Consequently there are also rather short intermolecular C...C distances (3.25, 3.36 Å).

The glide planes (010) at  $y = \pm \frac{1}{4}$  (Fig. 4) and the planes (001) at  $z \approx 0$  and  $z \approx \frac{1}{2}$  (Fig. 5) have cleavage character, because they are only penetrated by intermolecular distances. The layer character also facilitates the twinning found after cutting the crystals parallel to [010], because only a few intermolecular contacts approximately parallel to [100] are involved in the reorientation of the layers with a change of angle  $\beta$  into  $180^\circ - \beta$ . Finally the layer-type structure is in accordance with the observed orientation of the optical indicatrix, since  $n_y$  is parallel to the longest direction of the molecule and  $n_x$  is approximately perpendicular to the plane of the molecule.

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